

**REPRESENTATION OF SOLUTIONS TO LINEAR
ALGEBRAIC EQUATIONS**

by

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of a dissertation submitted by

Fred Krylov

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I have read the dissertation of Fred Krylov in its final form and have found that (1) its format, citations, and bibliographic style are consistent and acceptable; (2) its illustrative materials including figures, tables, and charts are in place; and (3) the final manuscript is satisfactory to the Supervisory Committee and is ready for submission to The Graduate School.

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ABSTRACT

This dissertation is described here in less than 350 words using no footnotes, diagrams, references or outside anything.

For my cat, Mouse, a few lines only.

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NOTATION AND SYMBOLS

Most of the following may be found in [4].

\mathbb{R}^n	n -dimensional Euclidean space.
Ω	A bounded open subset of \mathbb{R}^n .
$\partial\Omega$	The boundary of Ω .
B_r	$= \{x : x < r\}$, the ball of radius r .
$D^\beta f$	$= \partial^{ \beta } f / \partial x_{\beta_1} \partial x_{\beta_2} \cdots \partial x_{\beta_n}$, $ \beta \equiv \sum_i \beta_i$.
∇f	$= (\partial f / \partial x_1, \partial f / \partial x_2, \dots, \partial f / \partial x_n)$, the gradient of f .
$\operatorname{div}\{\mathbf{g}\}$	$= \partial g_1 / \partial x_1 + \partial g_2 / \partial x_2 + \dots + \partial g_n / \partial x_n$, the divergence of \mathbf{g} .
Δf	$= \operatorname{div}\{\nabla f\}$, the Laplacian of f .
$C^k(\Omega)$	Functions defined on Ω which have k continuous derivatives.
$C_0^k(\Omega)$	$C^k(\Omega)$ functions which vanish at the boundary.
$C^{0,\alpha}(\Omega)$	Hölder continuous functions with Hölder constant α .
$C^{k,\alpha}(\Omega)$	$C^k(\Omega)$ functions with $C^{0,\alpha}(\Omega)$ derivatives (up to order k).
$C_0^{k,\alpha}(\Omega)$	$C^{k,\alpha}(\Omega)$ functions which vanish at the boundary.
$\ f\ _{L^p(\Omega)}$	$= \left(\int_\Omega f ^p dx\right)^{1/p}$, the L^p norm.
$L^p(\Omega)$	The space of p integrable functions (the L^p norm is bounded).
$\ f\ _{W^{k,p}(\Omega)}$	$= \left(\sum_{ \beta \leq k} \int_\Omega D^\beta f ^p dx\right)^{1/p}$, the Sobolev norm.
$W^{k,p}(\Omega)$	The space of functions with bounded Sobolev norm.
$W_0^{k,p}(\Omega)$	$W^{k,p}(\Omega)$ functions that vanish a.e. at the boundary.

ACKNOWLEDGEMENTS

This page is optional. It's in the table of contents and it's labeled 'ACKNOWLEDGMENTS' even though the spelling 'ACKNOWLEDGEMENTS' is also correct. This page should be at the end of the preface, if one exists, or a separate page, if no preface is used.

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CHAPTER 1

INTRODUCTION

1.1 The Sample.tex file

1.1.1 Ordinary Text

The ends of words and sentences are marked by spaces. It doesn't matter how many spaces you type; one is as good as 100. The end of a line counts as a space.¹

One or more blank lines denote the end of a paragraph.

Since any number of consecutive spaces are treated like a single one, the formatting of the input file makes no difference to L^AT_EX, but it makes a difference to you. When you use L^AT_EX, making your input file as easy to read as possible will be a great help as you write your document and when you change it. This sample file shows how you can add comments to your own input file.

Because printing is different from typewriting, there are a number of things that you have to do differently when preparing an input file than if you were just typing the document directly. Quotation marks like “this” have to be handled specially, as do quotes within quotes: “‘this’ is what I just wrote, not ‘that’.”

Dashes come in three sizes: an intra-word dash, a medium dash for number ranges like 1–2, and a punctuation dash—like this.

A sentence-ending space should be larger than the space between words within a sentence. You sometimes have to type special commands in conjunction with punctuation characters to get this right, as in the following sentence. Gnats, gnus, etc. all begin with G. You should check the spaces after periods when reading your output to make sure you haven't forgotten any special cases. Generating an ellipsis . . . with the right spacing around the periods requires a special command.

L^AT_EX interprets some common characters as commands, so you must type special commands to generate them. These characters include the following: \$ & % # { and }.

¹This is a sample input file. Comparing it with the output it generates can show you how to produce a simple document of your own.

In printing, text is emphasized by using an *italic* type style.

A long segment of text can also be emphasized in this way. Text within such a segment given additional emphasis with Roman type. Italic type loses its ability to emphasize and become simply distracting when used excessively.

It is sometimes necessary to prevent L^AT_EX from breaking a line where it might otherwise do so. This may be at a space, as between the “Mr.” and “Jones” in “Mr. Jones,” or within a word—especially when the word is a symbol like *itemnum* that makes little sense when hyphenated across lines.

Footnotes² pose no problem.

L^AT_EX is good at typesetting mathematical formulas like $x - 3y = 7$ or

$$a_1 > x^{2n}/y^{2n} > x'.$$

Remember that a letter like x is a formula when it denotes a mathematical symbol, and should be treated as one.

1.1.2 Displayed Text

Text is displayed by indenting it from the left margin.

1.1.2.1 Quotations

Quotations are commonly displayed. There are short quotations

This is a short a quotation. It consists of a single paragraph of text. There is no paragraph indentation.

and longer ones.

This is a longer quotation. It consists of two paragraphs of text. The beginning of each paragraph is indicated by an extra indentation.

This is the second paragraph of the quotation. It is just as dull as the first paragraph.

1.1.2.2 Lists

Another frequently-displayed structure is a list.

²This is an example of a footnote.

1.1.2.2.1 Itemize. The following is an example of an *itemized* list.

- This is the first item of an itemized list. Each item in the list is marked with a “tick”. The document style determines what kind of tick mark is used.
- This is the second item of the list. It contains another list nested inside it. The inner list is an *enumerated* list.
 1. This is the first item of an enumerated list that is nested within the itemized list.
 2. This is the second item of the inner list. \LaTeX allows you to nest lists deeper than you really should.

This is the rest of the second item of the outer list. It is no more interesting than any other part of the item.

- This is the third item of the list.

1.1.2.2.2 Verse. You can even display poetry.

There is an environment for verse

Whose features some poets will curse.

For instead of making

Them do *all* line breaking,

It allows them to put too many words on a line when they'd
rather be forced to be terse.

1.1.2.3 Mathematics

Mathematical formulas may also be displayed. A displayed formula is one-line long; multiline formulas require special formatting instructions.

$$x' + y^2 = z_i^2$$

Don't start a paragraph with a displayed equation, nor make one a paragraph by itself.

1.2 More examples: Jeff McGough's Thesis

Equations like $\gamma = 0$ that don't need numbering may be set inline by the coding `$\gamma = 0$` or displayed by

\$\$
\gamma = 0.
\$\$

Numbered equations are set as shown in the next paragraph. They use the theorem environments defined in `thesis.sty`:

```
\newtheorem{thrm}{Theorem}
\newtheorem{lem}[thrm]{Lemma}
\newtheorem{cor}[thrm]{Corollary}
\newtheorem{rem}[thrm]{Remark}
\newtheorem{defn}[thrm]{Definition}
\newtheorem{exmpl}[thrm]{Example}
```

The Gelfand problem is the following elliptic boundary value problem:

$$\begin{aligned}\Delta u + \lambda e^u &= 0, & u \in \Omega, \\ u &= 0, & u \in \partial\Omega.\end{aligned}\tag{1.1}$$

The previous equation had a label. It may be referenced as equation (1.1).

1.3 History of the Gelfand problem

According to Bebernes and Eberly [1, p.46], Gelfand was “the first to make an in-depth study” of (1.1). Following this statement they briefly outline the history of the Gelfand problem.

For dimension $n = 1$, Liouville [6] first studied and found an explicit solution in 1853. For $n = 2$, Bratu [2] found an explicit solution in 1914. Frank-Kamenetski [3] rediscovered these results in his development of thermal explosion theory. Joseph and Lundgren [5] gave an elementary proof via phase plane analysis of the multiple existence of solutions for dimensions $n \geq 3$.

From Zeidler [8]:

$$\begin{aligned}\operatorname{div} j &= f, & x \in \Omega, \\ u &= g_1, & x \in \partial\Omega_1, \\ j\nu &= g_2, & x \in \partial\Omega_2,\end{aligned}\tag{1.2}$$

where

$$j = h(|\nabla u|^2)\nabla u\tag{1.3}$$

and Ω is a bounded domain in \mathbb{R}^n with smooth boundary $\partial\Omega = \overline{\partial\Omega_1} \cup \overline{\partial\Omega_2}$, $\partial\Omega_1 \cap \partial\Omega_2 = \emptyset$ and ν is the normal vector to $\partial\Omega$.

Lemma 1 *Assuming that $\partial\Omega_2 = \emptyset$ and that $h(t) = 1$, we have*

$$\begin{aligned}\Delta u &= f, & x \in \Omega, \\ u &= g_1, & x \in \partial\Omega.\end{aligned}$$

Corollary 2 *If $g_2 = 0$ then*

$$\begin{aligned}\Delta u &= f, & x \in \Omega, \\ u &= 0, & x \in \partial\Omega.\end{aligned}$$

1.4 Fundamental results

The investigation of the Gelfand problem begins with examining the (this paragraph continues for many lines).

Theorem 3 (Joseph-Lundgren [5]) *Boundary value problem (1.1) has positive radial solutions u on the unit ball which depend on n and λ in the following manner.*

1. *For $n = 1, 2$, there exists $\lambda^* > 0$ such that*
 - (a) *for $0 < \lambda < \lambda^*$ there are two positive solutions,*
 - (b) *for $\lambda = \lambda^*$ there is a unique solution, and*
 - (c) *for $\lambda > \lambda^*$ there are no solutions.*
2. *For $3 \leq n \leq 9$, let $\bar{\lambda} = 2(n - 2)$; then there exist positive constants λ_* , λ^* with $0 < \lambda_* < \bar{\lambda} < \lambda^*$, such that*
 - (a) *for $\lambda = \lambda^*$ there is a unique solution,*
 - (b) *for $\lambda > \lambda^*$ there are no solutions,*
 - (c) *for $\lambda = \bar{\lambda}$ there is a countably infinite number of solutions,*
 - (d) *for $\lambda \in (\lambda_*, \lambda^*)$, $\lambda \neq \bar{\lambda}$, there is a finite number of solutions,*
 - (e) *for $\lambda < \lambda_*$ there is a unique solution.*
3. *For $n \geq 10$, let $\lambda^* = 2(n - 2)$ then*
 - (a) *for $\lambda \geq \lambda^*$ there are no solutions,*
 - (b) *for $\lambda \in (0, \lambda^*)$ there is a unique solution.*

CHAPTER 2

QUADRATIC NONLINEARITIES

In this chapter we derive results for the quadratic equation.

2.1 Derivation of the quadratic formula

A quadratic equation is one of the form

$$ax^2 + bx + c = 0 \tag{2.1}$$

where a, b, c are known constants and x is the unknown. The results are summarized in Table 2.1 and Table 2.2 below.

2.2 Application of the quadratic formula

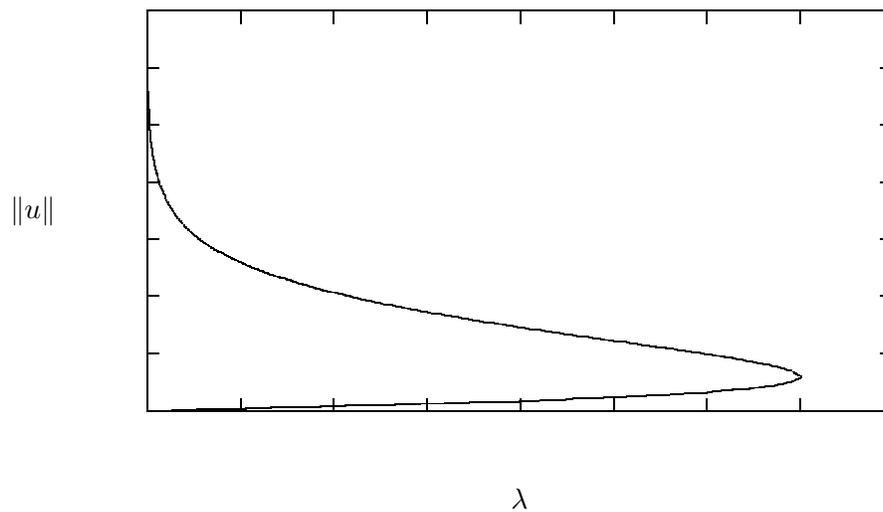
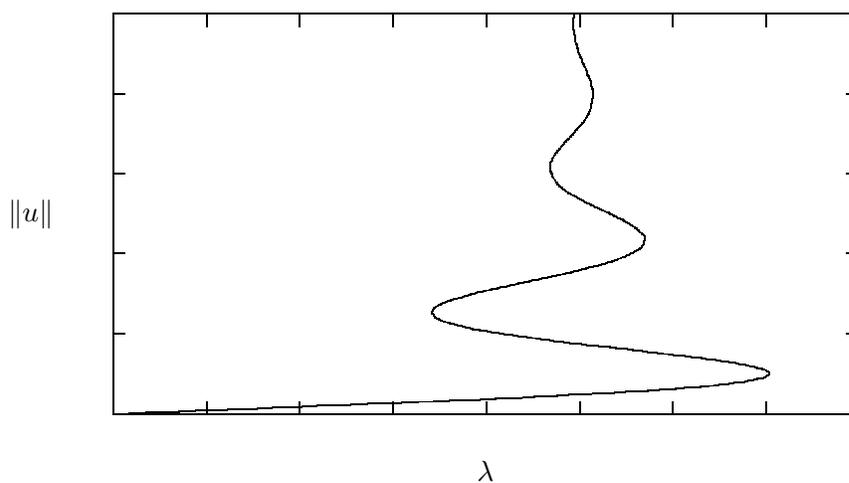
If the differential operator generates a nonnegative form, then an inequality is based on the following considerations. See Figure 2.1 for $n = 1, 2$, Figure 2.2 for $3 \leq n \leq 9$ and Figure 2.3 for $n \geq 10$.

Table 2.1. PDE solve times, $15^3 + 1$ equations.

Precond.	Time	Nonlinear Iterations	Krylov Iterations	Function calls	Precond. solves
None	1260.9u (21:09)	3	26	30	0
FFT	983.4u (16:31)	2	5	8	7
MILU	629.7u (10:36)	3	11	15	14

Table 2.2. Convergence properties of RQL.

Object	Normal Matrices	Diagonalizable Matrices	Defective Matrices
ρ	Stationary at ev's.	Stationary at ev's.	Stationary at ev's.
$\ r_k\ $	$\rightarrow 0$ as $k \rightarrow \infty$.	Can oscillate.	Can oscillate.
ρ_k	Converges.	Unknown.	Unknown.
Convergence to eigensets	is cubic.	is quadratic.	is linear.

**Figure 2.1.** Gelfand equation on the ball, $n = 1, 2$.**Figure 2.2.** Gelfand equation on the ball, $3 \leq n \leq 9$.

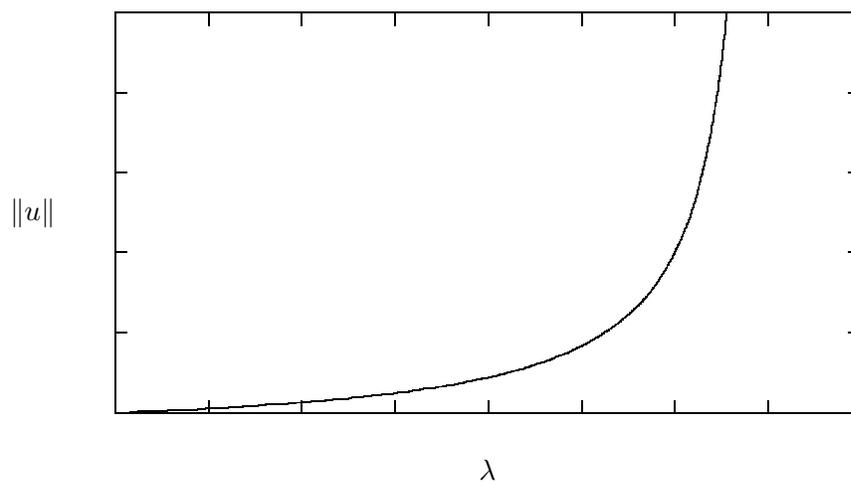


Figure 2.3. Gelfand equation on the ball, $n \geq 10$.

CHAPTER 3

SYSTEMS

3.1 Diagrams made with `diagram.sty`

An example diagram appears below in Figure 3.1. This is typical of what can be made with the `diagram` package.

3.2 Sample diagrams from `diagram.tex`

Example diagrams reproduced here were taken from various sources. Compare the three diagrams of increasing sizes in Figure 3.2, Figure 3.3, Figure 3.4 with the three diagrams in Figure 3.5, Figure 3.6, Figure 3.7.

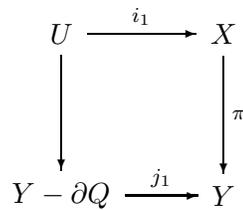


Figure 3.1. Diagram example

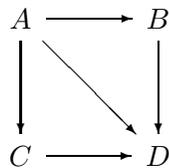


Figure 3.2. Base diagram, Arrowlength = 3.0em

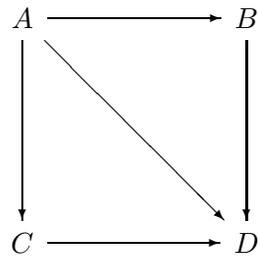


Figure 3.3. Same as Figure 3.2, but Arrowlength = 6.0em

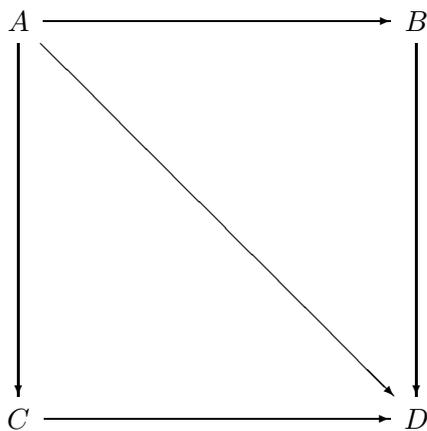


Figure 3.4. Same as Figure 3.2, but Arrowlength = 12.0em

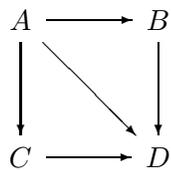


Figure 3.5. Base figure, same as Figure 3.2.

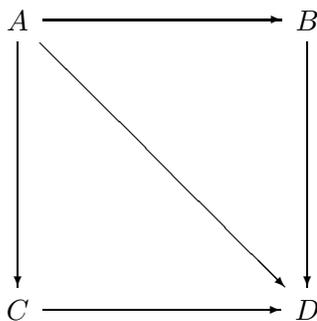


Figure 3.6. Same as Figure 3.5, but Bignode = strut hspace 6.0em.

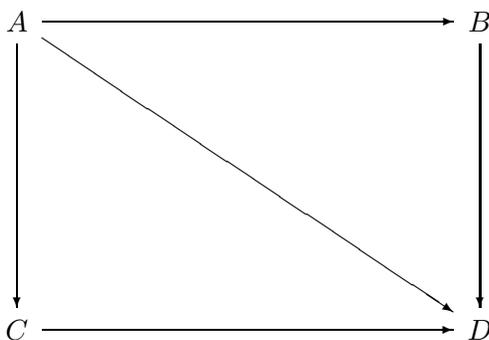


Figure 3.7. Same as Figure 3.2, but Bignode = strut hspace 12.0em

Below we show diagrams from the manual with a few modifications. The first in Figure 3.8 is essentially as it appears in the manual, whereas the second, Figure 3.9 has been rescaled to a larger size.

Below are several diagrams created by Bill Richter. The first, Figure 3.10 is modified slightly to produce Figure 3.11. Both use fractur fonts. The last one, Figure 3.12, is a complicated example illustrating the limits of what can be done with diagrams.

The diagram below in Figure 3.13, the last of our series of illustrations, is by Anders Thorup (thorup@math.ku.dk), originally done with a package developed by himself and Steven Kleiman (kleiman@math.mit.edu):

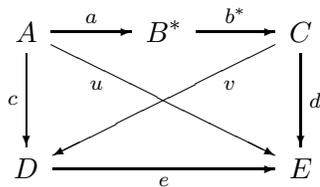


Figure 3.8. First diagram from manual

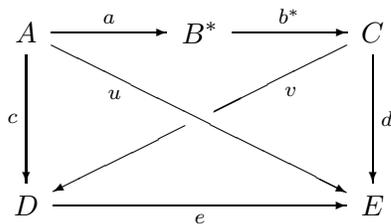


Figure 3.9. First diagram from manual, rescaled.

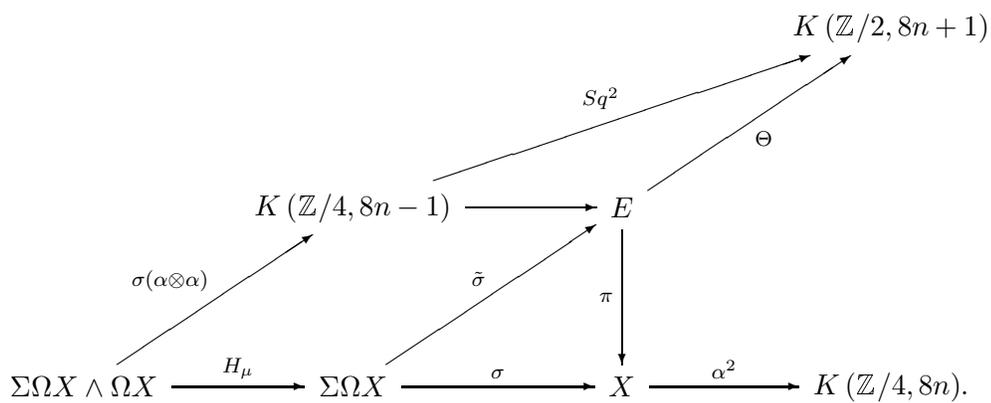


Figure 3.10. Bill Richter, first diagram

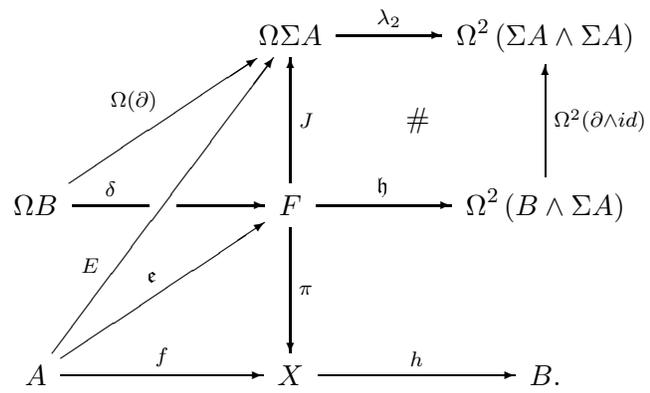


Figure 3.11. Bill Richter, second diagram

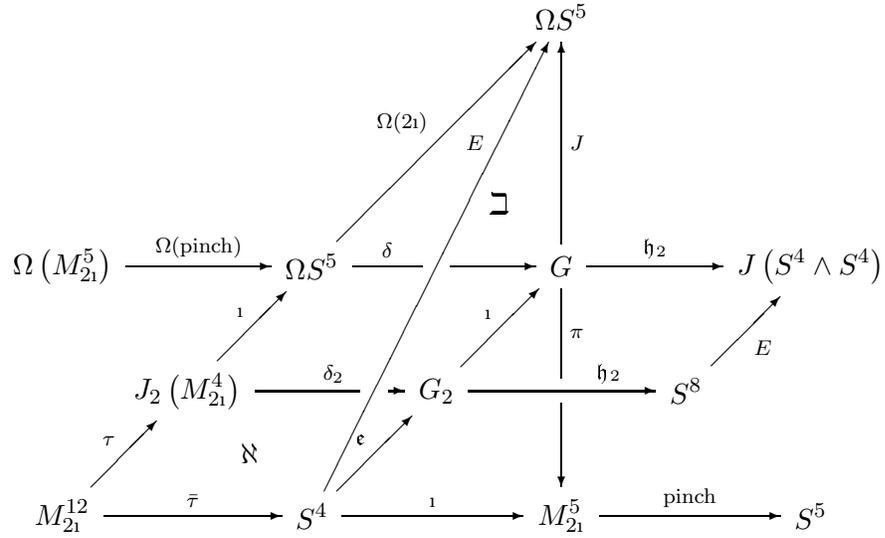


Figure 3.12. Bill Richter, third diagram

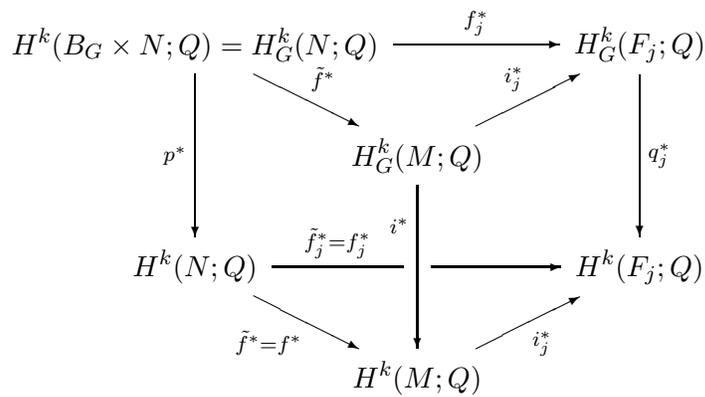


Figure 3.13. Anders Thorup diagram

APPENDIX

CLASSICAL IDENTITIES

Rellich's identity

Standard developments of Pohozaev's identity used an identity by Rellich [7], reproduced here.

Lemma 1 (Rellich) *Given L in divergence form and a, d defined above, $u \in C^2(\Omega)$, we have*

$$\begin{aligned} \int_{\Omega} (-Lu) \nabla u \cdot (x - \bar{x}) \, dx &= (1 - \frac{n}{2}) \int_{\Omega} a(\nabla u, \nabla u) \, dx - \frac{1}{2} \int_{\Omega} d(\nabla u, \nabla u) \, dx \\ &+ \frac{1}{2} \int_{\partial\Omega} a(\nabla u, \nabla u)(x - \bar{x}) \cdot \nu \, dS - \int_{\partial\Omega} a(\nabla u, \nu) \nabla u \cdot (x - \bar{x}) \, dS. \end{aligned} \quad (\text{A.1})$$

Proof:

There is no loss in generality to take $\bar{x} = 0$. First rewrite L :

$$Lu = \frac{1}{2} \left[\sum_i \sum_j \frac{\partial}{\partial x_i} \left(a_{ij} \frac{\partial u}{\partial x_j} \right) + \sum_i \sum_j \frac{\partial}{\partial x_i} \left(a_{ij} \frac{\partial u}{\partial x_j} \right) \right]$$

Switching the order of summation on the second term and relabeling subscripts, $j \rightarrow i$ and $i \rightarrow j$, then using the fact that $a_{ij}(x)$ is a symmetric matrix, gives the symmetric form needed to derive Rellich's identity.

$$Lu = \frac{1}{2} \sum_{i,j} \left[\frac{\partial}{\partial x_i} \left(a_{ij} \frac{\partial u}{\partial x_j} \right) + \frac{\partial}{\partial x_j} \left(a_{ij} \frac{\partial u}{\partial x_i} \right) \right]. \quad (\text{A.2})$$

Multiplying $-Lu$ by $\frac{\partial u}{\partial x_k} x_k$ and integrating over Ω , yields

$$\int_{\Omega} (-Lu) \frac{\partial u}{\partial x_k} x_k \, dx = -\frac{1}{2} \int_{\Omega} \sum_{i,j} \left[\frac{\partial}{\partial x_i} \left(a_{ij} \frac{\partial u}{\partial x_j} \right) + \frac{\partial}{\partial x_j} \left(a_{ij} \frac{\partial u}{\partial x_i} \right) \right] \frac{\partial u}{\partial x_k} x_k \, dx$$

Integrating by parts (for integral theorems see [8, p. 20]) gives

$$= \frac{1}{2} \int_{\Omega} \sum_{i,j} a_{ij} \left[\frac{\partial u}{\partial x_j} \frac{\partial^2 u}{\partial x_k \partial x_i} + \frac{\partial u}{\partial x_i} \frac{\partial^2 u}{\partial x_k \partial x_j} \right] x_k \, dx$$

$$\begin{aligned}
& + \frac{1}{2} \int_{\Omega} \sum_{i,j} a_{ij} \left[\frac{\partial u}{\partial x_j} \delta_{ik} + \frac{\partial u}{\partial x_i} \delta_{jk} \right] \frac{\partial u}{\partial x_k} dx \\
& - \frac{1}{2} \int_{\partial\Omega} \sum_{i,j} a_{ij} \left[\frac{\partial u}{\partial x_j} \nu_i + \frac{\partial u}{\partial x_i} \nu_j \right] \frac{\partial u}{\partial x_k} x_k dx
\end{aligned}$$

= $I_1 + I_2 + I_3$, where the unit normal vector is ν . One may rewrite I_1 as

$$I_1 = \frac{1}{2} \int_{\Omega} \sum_{i,j} a_{ij} \frac{\partial}{\partial x_k} \left(\frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} \right) x_k dx$$

Integrating the first term by parts again yields

$$\begin{aligned}
I_1 &= -\frac{1}{2} \int_{\Omega} \sum_{i,j} a_{ij} \left(\frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} \right) dx + \frac{1}{2} \int_{\partial\Omega} \sum_{i,j} a_{ij} \left(\frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} \right) x_k \nu_k dS \\
&\quad - \frac{1}{2} \int_{\Omega} \sum_{i,j} \left(\frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} \right) x_k \frac{\partial a_{ij}}{\partial x_k} dx.
\end{aligned}$$

Summing over k gives

$$\begin{aligned}
& \int_{\Omega} (-Lu)(\nabla u \cdot x) dx = -\frac{n}{2} \int_{\Omega} \sum_{i,j} a_{ij} \left(\frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} \right) dx \\
& + \frac{1}{2} \int_{\partial\Omega} \sum_{i,j} a_{ij} \left(\frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} \right) (x \cdot \nu) dS - \frac{1}{2} \int_{\Omega} \sum_{i,j} \left(\frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} \right) (x \cdot \nabla a_{ij}) dx \\
& + \frac{1}{2} \int_{\Omega} \sum_{i,j,k} a_{ij} \left[\frac{\partial u}{\partial x_j} \frac{\partial u}{\partial x_k} \delta_{ik} + \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_k} \delta_{jk} \right] dx \\
& - \frac{1}{2} \int_{\partial\Omega} \sum_{i,j} a_{ij} \left[\frac{\partial u}{\partial x_j} \nu_i + \frac{\partial u}{\partial x_i} \nu_j \right] (\nabla u \cdot x) dS.
\end{aligned}$$

Combining the first and fourth term on the right-hand side simplifies the expression

$$\begin{aligned}
& \int_{\Omega} (-Lu)(\nabla u \cdot x) dx = \left(1 - \frac{n}{2}\right) \int_{\Omega} \sum_{i,j} a_{ij} \left(\frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} \right) dx \\
& + \frac{1}{2} \int_{\partial\Omega} \sum_{i,j} a_{ij} \left(\frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} \right) (x \cdot \nu) dS - \frac{1}{2} \int_{\Omega} \sum_{i,j} \left(\frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} \right) (x \cdot \nabla a_{ij}) dx \\
& - \frac{1}{2} \int_{\partial\Omega} \sum_{i,j} a_{ij} \left[\frac{\partial u}{\partial x_j} \nu_i + \frac{\partial u}{\partial x_i} \nu_j \right] (\nabla u \cdot x) dS.
\end{aligned}$$

Using the notation defined above, the result follows.

Fortran code

```

subroutine rhs(neq,v,rhsf)
save
c
c This subroutine computes the function values. Inputs are neq and
c v, and on output the values of f are stored in the array of rhsf
c
include at10.0pt'parabolic.inc'

integer neq
integer i
integer j
integer k
integer ind
integer inde
integer indw
integer indn
integer inds
integer ind0
integer ind1
integer ind2

double precision v(neq)
double precision rhsf(neq)
double precision u(nv)
double precision diff
double precision diffn
double precision diffxn
double precision diffyn
double precision nl

c write(*,*)'funct begin'

c
c Compute F for the local dynamics, written as  $F(u) = -du/dt + f(u)$ 
c
c
c the system parameters
c
c p1 ! parameter F
c p2 ! parameter k

do j = 1, ny
do i = 1, nx

c
c set up index
c
ind = (i-1)*nv + (j-1)*neq
c
c Extract the jth component at current time
c

```

```

nl = v(1+ind)*v(2+ind)*v(2+ind)

rhsf(1+ind) = (- nl + p1*(1.0d0 - v(1+ind)))*local
rhsf(2+ind) = ( nl - (p1+p2)*v(2+ind))*local

      end do
    end do

c -----
c
c   add diffusion for all species (zero diffusion
c   coefficient takes care of those that do not diffuse).
c
c -----

      do j = 1, ny
        do i = 1, nx

c indexing
c
          ind0 = (i-1)*nv + (j-1)*meq  ! point
          indw = (i-2)*nv + (j-1)*meq  ! west point
          inde = (i)*nv + (j-1)*meq    ! east point
          indn = (i-1)*nv + (j)*meq    ! north point
          inds = (i-1)*nv + (j-2)*meq  ! south point

          if(i.eq.1) indw = (nx-1)*nv + (j-1)*meq
          if(i.eq.nx) inde = (j-1)*meq
          if(j.eq.1) inds = (i-1)*nv + (ny-1)*meq
          if(j.eq.ny) indn = (i-1)*nv

          do k = 1, 2

c
c   First compute the contribution within a row at the current time
c   and at the preceding time.
c
            ind = k + ind0
            ind1 = k + indw
            ind2 = k + inde

            diffxn = v(ind1) - 2.0d0*v(ind) + v(ind2)

c
c   Compute the contribution from the columns
c
            ind1 = k + indn
            ind2 = k + inds

            diffyn = v(ind1) - 2.0d0*v(ind) + v(ind2)

c

```

c Multiply by other factors and sum

c

$$\text{diff} = d(k) * hxx * (\text{diffxn} + \text{diffyn}) * \text{diffus}$$

$$\text{rhsf(ind)} = \text{rhsf(ind)} + \text{diff}$$

110

end do

end do

end do

return

end

120

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